Reconstruction of Neutron Cross-sections and Sampling

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Outline

- ☐ Introduction
- Reconstruction of resonance cross-section
- Linearization of cross-section
- Unionization
- ☐ Doppler broadening at higher temperature
- Sampling
- ☐ Independent angular distribution
- ☐ Independent energy distribution
- Energy-angle correlated distribution
- ☐ Fission fragments and fission neutrons
- Summary



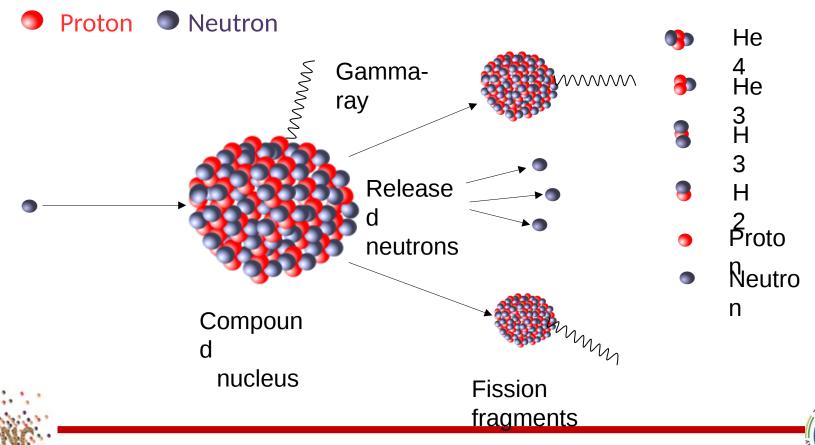
Introduction: Importance of low energy neutrons

- ➤ Most of the neutron applications are in low energy region (<20MeV) i.e. material studies/diffraction, fusion and fission reactors, Nuclear medicine, Radiation dosimetry in accelerator and nuclear devices etc.
- Low energy neutron transport takes significant time in hadron transport because of charge neutrality.
- Radiation dosimetry and shielding calculations in GEANT4 is not comparable with experimental data.



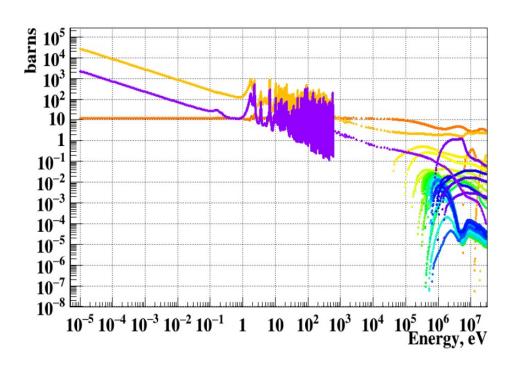
Introduction: Compound nucleus reactions

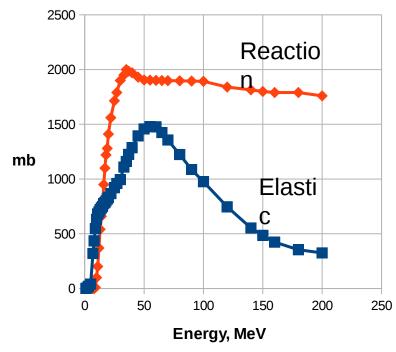
- The absence of coulomb barrier between neutron and nucleus as compared to charge particles makes neutron interactions special.
- It can penetrate deep inside the nucleus even at meV energies.



Introduction: neutron and proton cross-sections

Neutron interactions below 20MeV or 200MeV in some cases.





Neutron crosssections

Proton crosssections

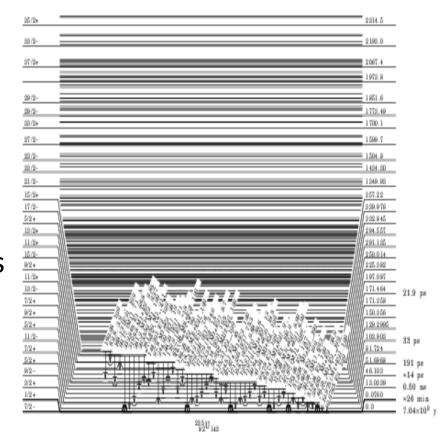


Introduction: Evaluated Nuclear Data File

- ➤ Whether model can predict the cross-sections? No
- ➤ Nuclear structure contribute to the final states.
- ➤ No single model for all the Isotopes that can work reasonably well
- ➤ What is the alternative solution?
- ► Evaluated Nuclear Data

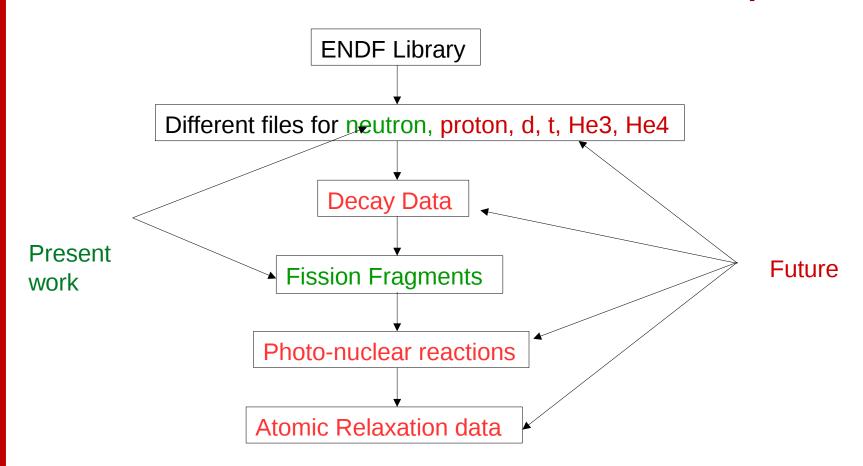
Disadvantage

Too many data points due to resonance



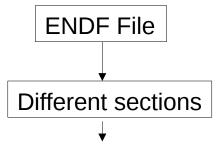


Introduction: Evaluated Nuclear Data Library





Introduction: Evaluated Nuclear Data File



- 1 → General description, Fission neutron multiplicity, partial photon data
- 2 → Resonance parameters for cross-section
- 3 → Cross-sections for all reactions
- 4 → Independent angular distributions
- 5 → Independent energy distributions
- 6 → Correlated angle-energy distributions
- 7 → Thermal neutron scattering data
- 8 → Decay data and fission products
- 9 → Multiplicities for radio-active nuclide production
- 10 → Production cross-section for radio-active nuclide
- 11 → General comments for Photon production
- 12 → Photon production multiplicities
- 13 → Photon production cross-section
- 14 → Photon angular distributions
- 15 → Continuous photon energy distribution
 - . 12 More sections about atomic reactions, errors, photon, election interaction

Introduction: Evaluated Nuclear Data File

How the data file look

```
[MAT, 2,151/ ZA, AWR, 0, 0, NIS, 0] HEAD (NIS=1)
[MAT, 2,151/ ZAI, ABN, 0,LFW, NER, 0]CONT (ZAI=ZA,ABN=1,LFW=0,NER=1)
[MAT, 2,151/ EL, EH, LRU, LRF, NRO, NAPS] CONT (LRU=0, LRF=0, NRO=0, NAPS=0)
[MAT, 2,151/ SPI, AP, 0, 0, NLS, 0]CONT (NLS=0)}
[MAT, 2, 0/0.0, 0.0, 0, 0, 0] SEND
[MAT, 0, 0/0.0, 0.0, 0, 0, 0] FEND
9.223500+4 2.330248+2
                                                        09228 2151 1
9.223500+4 1.000000+0
                                                        09228 2151 2
                                                        19228 2151 3
1.000000-5 2.250000+3
3.500000+0 9.602000-1
                                                        39228 2151 4
2.330200+2 9.602000-1
                                                      31939228 2151
                                          19158
2.038300+3 3.000000+0 1.970300-2 3.379200-2-4.665200-2-1.008800-19228 2151
```

There are many different sub-sections with different set of parameters and different structures

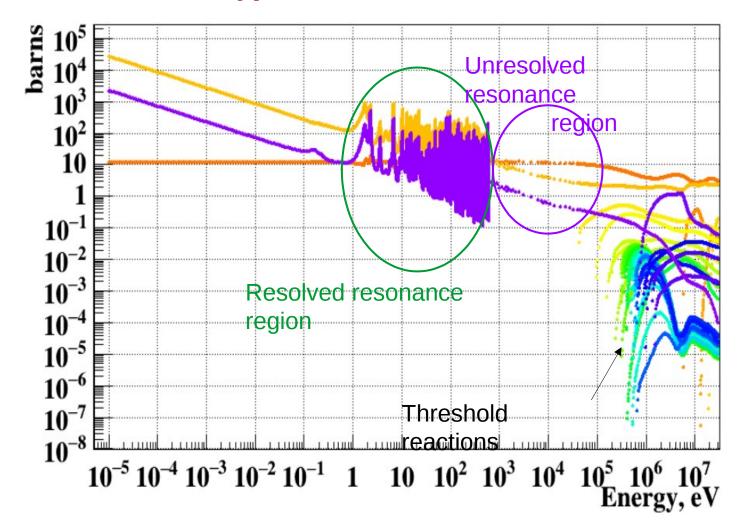


Introduction: Why point data are not given but parameters

- Too many data point
- > We don't understand from collection of points (lack of information about physics)
- ➤ Data should be interpreted by nuclear theory so that one can understand the physics
- ➤ We should be able to extrapolate and interpolate in the missing energy range
- Experimental information is utmost important to derive useful data.
- Further up-gradation of data is possible



Introduction: Typical cross-section







Reconstruction: Convert ASCII file to ROOT file

Read all sections from neutron data file

Read data from sub-libraries

- a) fission fragment yield
- b) decay data
- C)
- d)

.

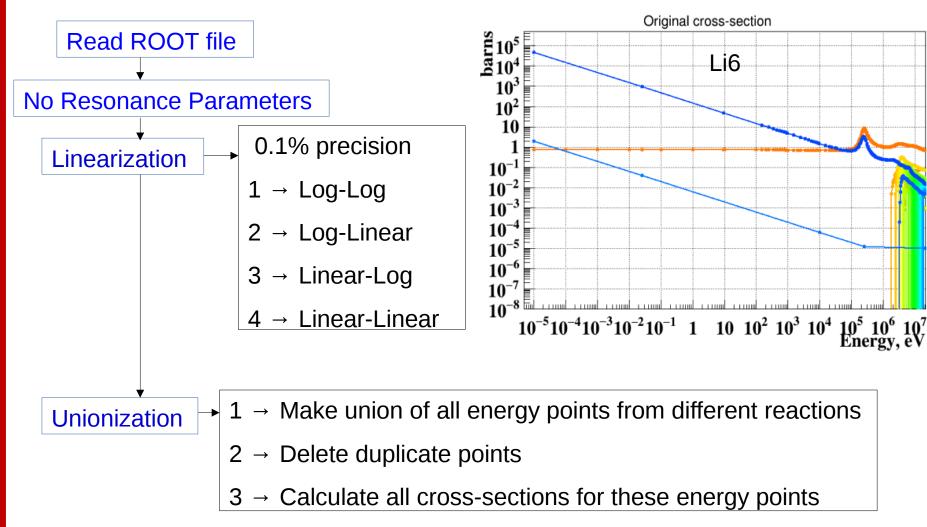
Convert 9.223500+4 data structure into Doubles/Float

Store into ROOT file → file size reduces 2-3 times

This is done before simulation into offline mode but one can do during simulation and go for a Coffee break



Reconstruction: Linearization and Unionization







Reconstruction: Total cross-section Total cross-section Read ROOT file No Resonance Parameters Tota Linearization Unionization Doppler broadening at higher temperature 10⁻⁵10⁻⁴10⁻³10⁻²10⁻¹ 1 10 10² 10³ 10⁴ 10⁵ 10⁶ 10⁷ Energy, eV Construct Total cross-section 1 → Add all cross-sections and generate total cross-section Modify ROOT file to replace

This is done before simulation into offline mode but one can do during simulation and go for a Lunch break

2 → Store only non zero data points



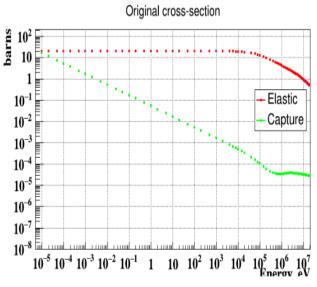


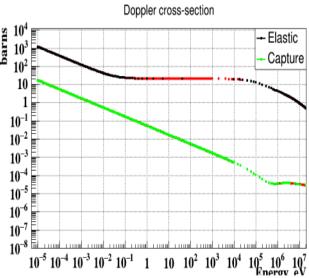
old cross-sections with new

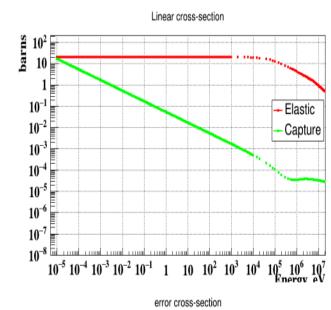
linearized point cross-sections

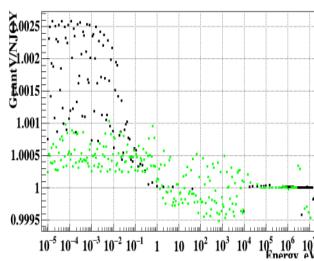
Reconstruction: Hydrogen cross-section

- Linear in log scale96 data points
- Linear data points487 data points
- Doppler broadning at 293.6Kelvin
- Linearization is
 to be done to
 gain some memory
 sometimes







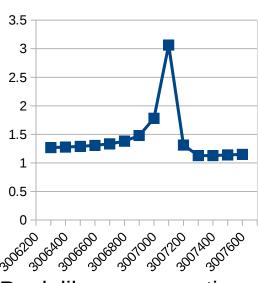




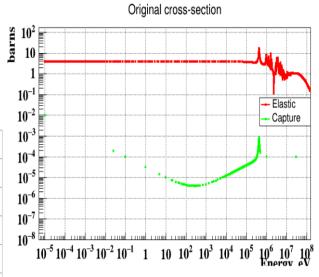


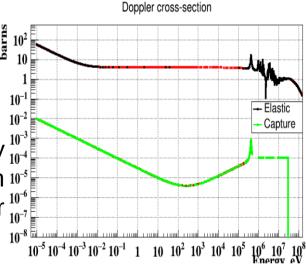
Reconstruction: 016 cross-section

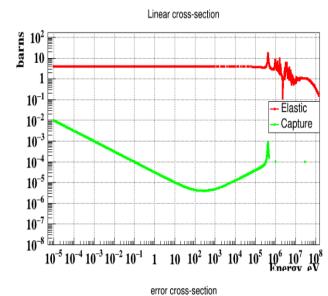
Data are given up to 200 MeV

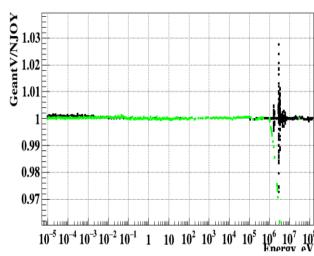


Peak like cross-section 10⁻²
But linearly 10⁻³
interpolatable Which 10⁻⁵
gives some Error 10⁻⁷
compared to NJOY 10⁻⁸













Reconstruction: Resonance cross-section

Read Resonance Parameters

1 → Calculate phase shifts, shift factors, penetration factors for higher angular momenta

1 → Resonance energy points

 $2 \rightarrow$ Resonance widths

3 → Resonance types

(Single level Breit-Wigner,

Multi-level Breit-Wigner,

Reich-Moore, Adler-Adler, R-Matrix)

Single level Breit-Wigner → 8 isotopes

Multi- level Breit-Wigner → 268 isotopes

Reice-Moore → 54 isotopes (best results)

Adler-Adler → None

 $R-Matrix \rightarrow None$ (Very hard to implement)



Reconstruction: Single Level Breit Wigner

$$\sigma_{n,n}(E) = \sum_{l=0}^{\text{NLS}-1} \sigma_{n,n}^l(E),$$

Elastic cross-section

$$\sigma_{n,n}^{l}(E) = (2l+1)\frac{4\pi}{k^{2}}\sin^{2}\phi_{l} + \frac{\pi}{k^{2}}\sum_{I}g_{J}\sum_{r=1}^{NR_{J}}\frac{\Gamma_{nr}^{2} - 2\Gamma_{nr}\Gamma_{r}\sin^{2}\phi_{l} + 2(E - E_{r}')\Gamma_{nr}\sin(2\phi_{l})}{(E - E_{r}')^{2} + \frac{1}{4}\Gamma_{r}^{2}}$$

Capture cross-section

$$\sigma_{n,\gamma}(E) = \sum_{l=0}^{\text{NLS}-1} \sigma_{n,\gamma}^l(E)$$

$$\sigma_{n,\gamma}^{l}(E) = \frac{\pi}{k^2} \sum_{J} g_J \sum_{r=1}^{NR_J} \frac{\Gamma_{nr} \Gamma_{\gamma r}}{(E - E_r')^2 + \frac{1}{4}\Gamma_r^2}$$

$$\sigma_{n,f}(E) = \sum_{l=0}^{\text{NLS}-1} \sigma_{n,f}^{l}(E) ,$$

$$\sigma_{n,f}^{l}(E) = \frac{\pi}{k^2} \sum_{J} g_{J} \sum_{r=1}^{NR_{J}} \frac{\Gamma_{nr} \Gamma_{fr}}{(E - E_{r}')^2 + \frac{1}{4} \Gamma_{r}^2}$$





Reconstruction: Multi-Level Breit Wigner

$$\sigma_{n,n}^{l(R)}(E) = \frac{\pi}{k^2} \sum_{J} g_J \sum_{r=1}^{NR_J} \frac{G_r \Gamma_r + 2H_r (E - E_r)}{(E - E_r')^2 + (\Gamma_r/2)^2}$$

Elastic cross-section

$$G_{r} = \frac{1}{2} \sum_{r'=1, r' \neq r}^{NR_{J}} \frac{\Gamma_{nr} \Gamma_{nr'} (\Gamma_{r} + \Gamma_{r'})}{(E'_{r} - E'_{r'})^{2} + \frac{1}{4} (\Gamma_{r} + \Gamma_{r'})^{2}},$$

$$H_{r} = \sum_{r'=1}^{NR_{J}} \frac{\Gamma_{nr} \Gamma_{nr'} (E_{r} - E_{r'})}{(E'_{r} - E'_{r'})^{2} + \frac{1}{4} (\Gamma_{r} + \Gamma_{r'})^{2}}$$

Capture cross-section

$$\sigma_{n,\gamma}(E) = \sum_{l=0}^{\text{NLS}-1} \sigma_{n,\gamma}^l(E)$$

$$\sigma_{n,\gamma}^{l}(E) = \frac{\pi}{k^2} \sum_{J} g_J \sum_{r=1}^{NR_J} \frac{\Gamma_{nr} \Gamma_{\gamma r}}{(E - E_r')^2 + \frac{1}{4}\Gamma_r^2}$$

$$\sigma_{n,f}(E) = \sum_{l=0}^{\text{NLS}-1} \sigma_{n,f}^{l}(E) ,$$

$$\sigma_{n,f}^{l}(E) = \frac{\pi}{k^2} \sum_{J} g_{J} \sum_{r=1}^{NR_{J}} \frac{\Gamma_{nr} \Gamma_{fr}}{(E - E_{r}')^2 + \frac{1}{4} \Gamma_{r}^2}$$



Reconstruction: Reich-Moore

Elastic cross-section

$$\sigma_{abs}(E) = \sigma_T(E) - \sigma_{nn}(E)$$

Absorption - fission

$$\sigma_T(E) = \frac{2\pi}{k^2} \sum_{l=0}^{\text{NLS}-1} \sum_{s=|I-\frac{1}{2}|}^{I+\frac{1}{2}} \sum_{J=|l-s|}^{l+s} g_J \text{Re} \left[1 - U_{lsJ,lsJ}\right]$$

$$\sigma_{nn}(E) = \frac{2\pi}{k^2} \sum_{l=0}^{\text{NLS}-1} \sum_{s=|I-\frac{1}{2}|}^{I+\frac{1}{2}} \sum_{J=|l-s|}^{l+s} g_J |1 - U_{lsJ,lsJ}|^2$$

$$\sigma_f(E) = \frac{2\pi}{k^2} \sum_{l=0}^{\text{NLS}-1} \sum_{s=|I-\frac{1}{2}|}^{I+\frac{1}{2}} \sum_{J=|l-s|}^{l+s} g_J \left[\left| U_{nf1}^{lsJ} \right|^2 + \left| U_{nf2}^{lsJ} \right|^2 \right]$$

$$U_{nb}^{J} = e^{-i(\phi_n + \phi_b)} \left\{ 2 \left[(I - K)^{-1} \right]_{nb} - \delta_{nb} \right\},\,$$

$$(I - K)_{nb} = \delta_{nb} - \frac{i}{2} \sum_{r} \frac{\Gamma_{nr}^{1/2} \Gamma_{br}^{1/2}}{E_r - E - i \Gamma_{\gamma r}/2}$$





Reconstruction: Unresolved resonance

$$\sigma_{n,n}(E) = \sum_{l=0}^{\text{NLS}-1} \sigma_{n,n}^{l}(E),$$

$$\sigma_{n,n}^{l}(E) = \frac{4\pi}{k^{2}} (2l+1) \sin^{2} \phi_{l}$$

$$+ \frac{2\pi^{2}}{k^{2}} \sum_{J}^{\text{NJS}} \left[\frac{g_{J}}{\overline{D}_{l,J}} \left\langle \frac{\Gamma_{n} \Gamma_{n}}{\Gamma} \right\rangle_{l,J} - 2\overline{\Gamma}_{nl,J} \sin^{2} \phi_{l} \right]$$

Elastic cross-section

Average widths are used along with fluctuation $\left\langle \frac{\Gamma_n \Gamma_n}{\Gamma} \right\rangle_{l,J} = \left(\frac{\overline{\Gamma}_{nl,J} \overline{\Gamma}_{nl,J}}{\overline{\Gamma}_{l,J}} \right) R_{n,l,J}$

Width fluctuation parameter R is calculated using MC²-II method.

Capture cross-section

$$\sigma_{n,\gamma}(E) = \sum_{l=0}^{\text{NLS}-1} \sigma_{n,\gamma}^{l}(E),$$

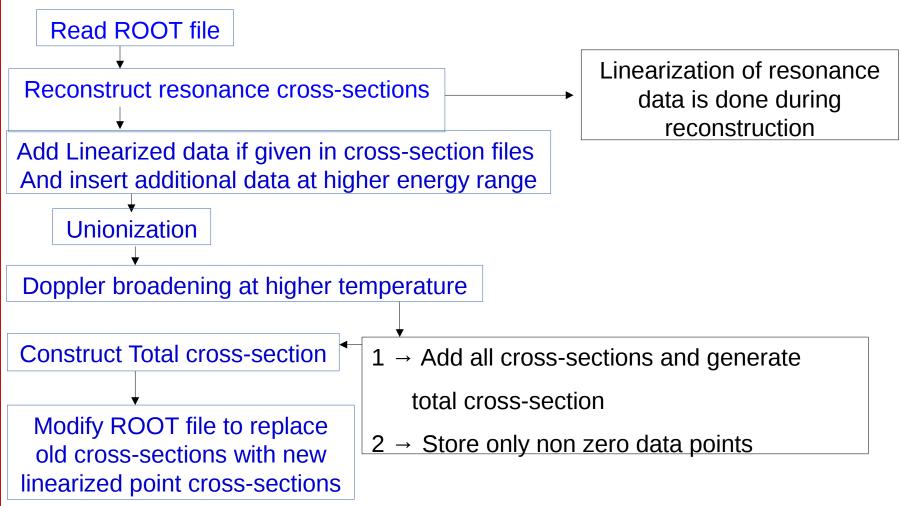
$$\sigma_{n,\gamma}^{l}(E) = \frac{2\pi^{2}}{k^{2}} \sum_{I}^{\text{NJS}} \frac{g_{J}}{\overline{D}_{l,J}} \left\langle \frac{\Gamma_{n} \Gamma_{\gamma}}{\Gamma} \right\rangle_{l,J}$$

$$\sigma_{n,f}(E) = \sum_{l=0}^{\text{NLS}-1} \sigma_{n,f}^{l}(E),$$

$$\sigma_{n,f}^{l}(E) = \frac{2\pi^{2}}{k^{2}} \sum_{J}^{\text{NJS}} \frac{g_{J}}{\overline{D}_{l,J}} \left\langle \frac{\Gamma_{n} \Gamma_{f}}{\Gamma} \right\rangle_{l,J}$$



Reconstruction: Resonance cross-section

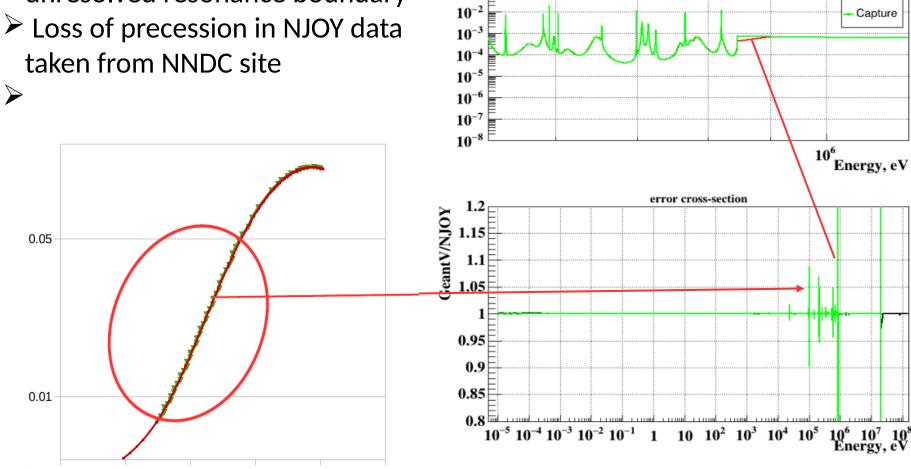


This is done before simulation into offline mode but one can do during simulation and go for a Lunch break



Reconstruction: Al27 cross-section

- Discontinuity at resolved and unresolved resonance boundary
- taken from NNDC site



10-



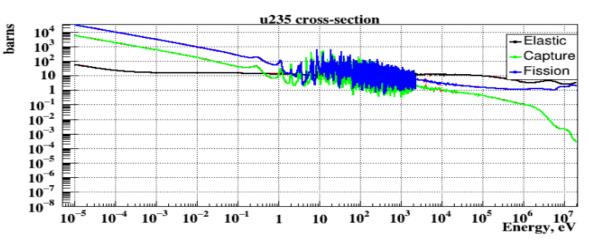
al27 cross-section

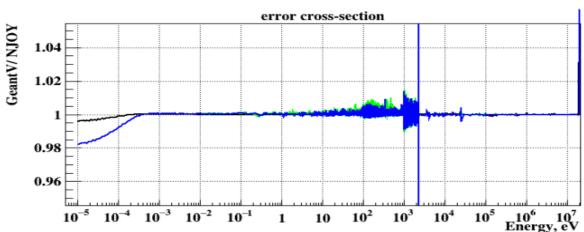
Reconstruction: U235 cross-sections

Data are given up to 30 MeV

Resolved and un-resolved resonance boundary shows discrepancy due to discontinuity

RR data are agreeing within 0.5%



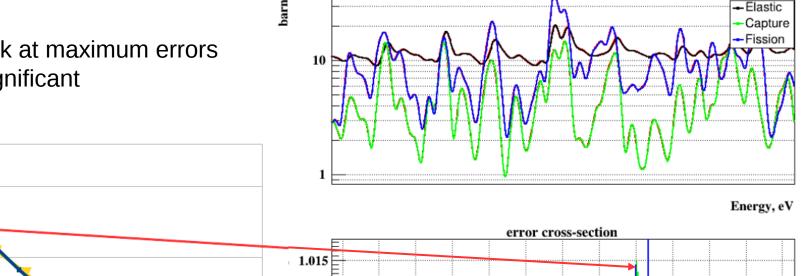


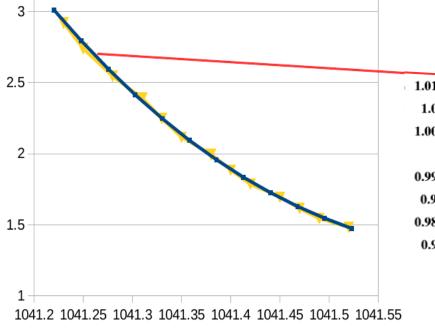
This is done before simulation into offline mode but one can do during simulation and go for a Day break

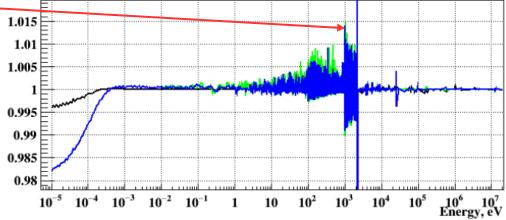


Reconstruction: U235 cross-sections

Closer look at maximum errors looks insignificant





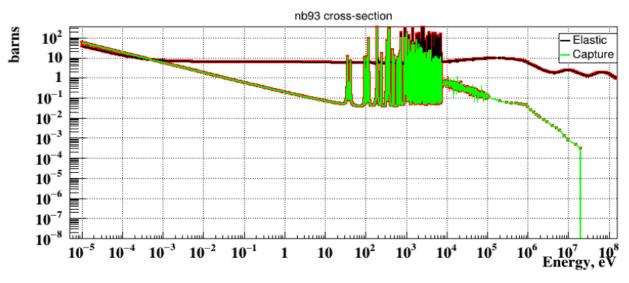


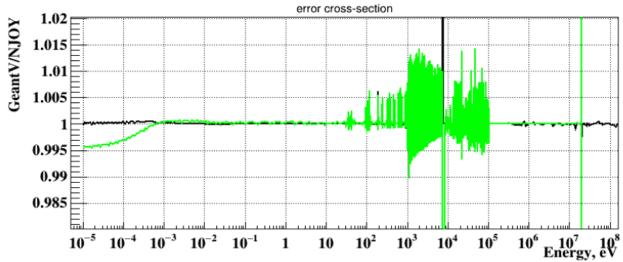
u235 cross-section





Reconstruction: Nb93 cross-sections









Reconstruction: Doppler broadening

Maxwellian velocity Distribution is used 0 Kelvin For the target atoms at 10^{-1} 1000 Kelvin different temperatures 10^{-2} 10^{-4} It is adopted from 10^{-5} **Federico's fortran** 10^{-6} version

Energy, eV





Reconstruction: Angular Distributions

- Angular distributions are given in terms of Legendre coefficients and probability tables
- ➤ Data are given mostly for few energies
- Cumulative distribution and PDF are used to get the angle

$$f(\mu,E) = \frac{2\pi}{\sigma_s(E)} \ \sigma(\mu,E) = \sum_{l=0}^{\rm NL} \frac{2l+1}{2} \ a_l(E) \ P_l(\mu)$$
 Legendre coefficients
$$E_i < E_{in} < E_{i+1}$$

$$E_{in} = E_i + r(E_{i+1} - E_i)$$

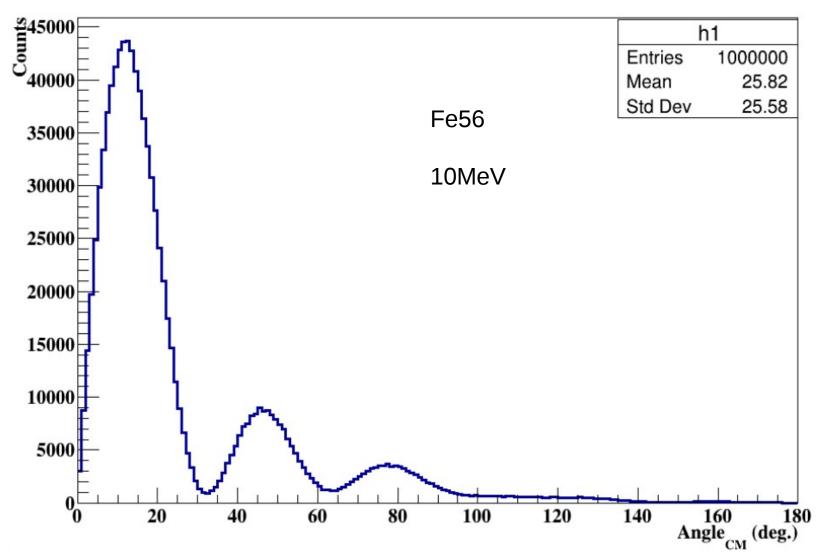
$$c_{l, k} < \xi_1 < c_{l, k+1}$$

$$\mu' = \mu_{l,k} + \left\{ \frac{\sqrt{P_{l,k}^2 + 2\left[\frac{p_{l,k+1} - p_{l,k}}{\mu_{l,k+1} - \mu_{l,k}}\right](\xi_1 - c_{l,k}) - p_{l,k}}}{\left[\frac{p_{l,k+1} - p_{l,k}}{\mu_{l,k+1} - \mu_{l,k}}\right]} \right\}$$

The making of probability tables are done at initialization and we plan to shift to offline Otherwise one can have a chat



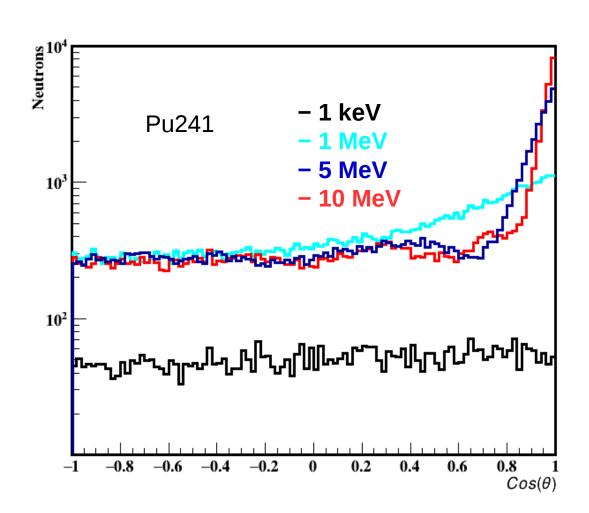
Reconstruction: Elastic Angular Distribution





Sampling: Elastic Angular Distribution

- ➤ Data are given mostly for few energies
- ➤ Bilinear interpolation
- ➤ Isotropic behavior for keV neutrons and forward peaking at higher energies
- ➤1 Million events are simulated





Reconstruction: Energy Distributions

- Energy distributions aregiven by tabular data or 5-6 different formulations
- We make all formats into probability tables
- Cumulative distribution and PDF are used to get the energy

One of the formulation for energy spectra

$$f(E \to E') = \frac{1}{2} [g(E', E_F(L)) + g(E', E_F(H))]$$

$$g(E', E_F) = \frac{1}{3\sqrt{(E_F T_M)}} \left[u_2^{3/2} E_1(u_2) - u_1^{3/2} E_1(u_1) + \gamma \left(\frac{3}{2}, u_2\right) - \gamma \left(\frac{3}{2}, u_1\right) \right]$$

$$u_1 = \left(\sqrt{E'} - \sqrt{E_F}\right)^2 / T_M$$

$$u_2 = \left(\sqrt{E'} + \sqrt{E_F}\right)^2 / T_M$$

 $E_F(X)$ are constant, which represent the average kinetic energy per nucleon of the fission fragment; arguments L and H refer to the average light fragment (given by the parameter EFL in the file) and the average heavy fragment (given by the parameter EFH in the file), respectively.

 T_M parameter tabulated as a function of incident neutron energy,

 $E_1(x)$ is the exponential integral,

 $\gamma(a,x)$ is the incomplete gamma function. The integral of this spectrum between zero and infinity is one. The value of the integral for a finite integration

Probability tables are made at initialization and we plan to shift to offline Otherwise one can have one more chat

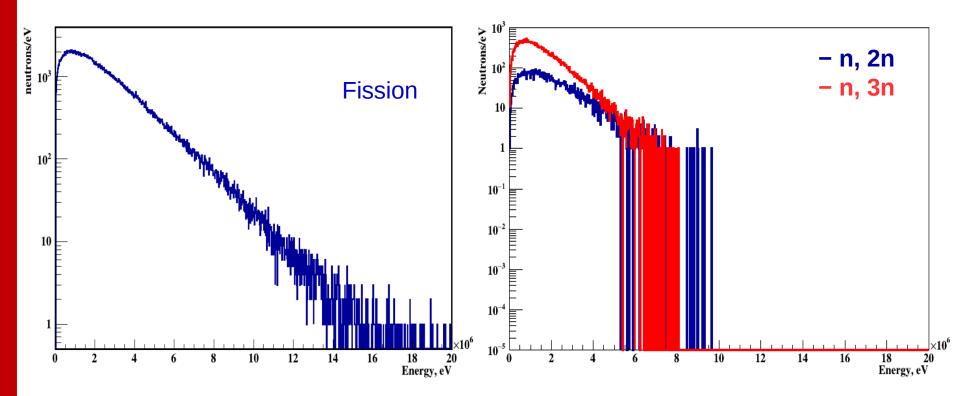


Sampling: Energy Distribution

In case of Fission second or third or higher number of neutrons are sampled from the same distribution

Second neutron from n,2n reaction can have Total - $\mathbf{1}^{st}$ neutron energy – recoil energy

Energy conservation exist but without correlation until such data are given

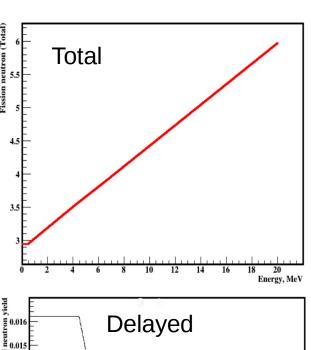


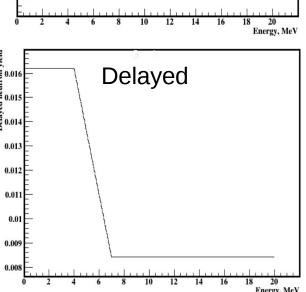


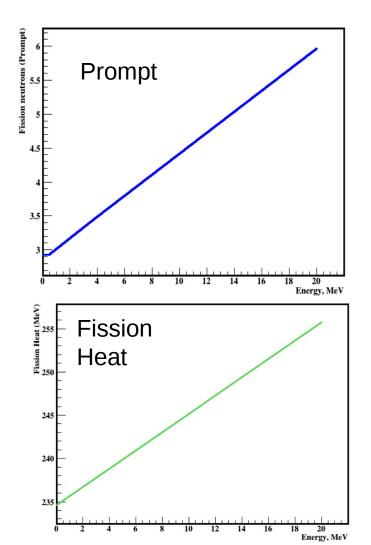


Sampling: PU-241 Fission neutrons

Sampling of average fission neutrons i.e. 2.56 is done based one Poissonian distribution





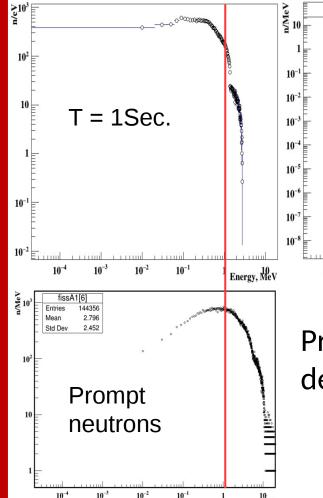


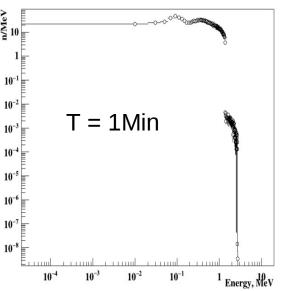


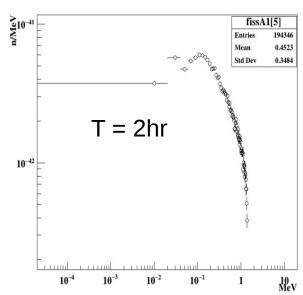


Sampling: PU-241 Delayed Fission neutrons

Delayed neutrons emitted by 6 precursor families. Mean time of decay is up to 100 seconds.







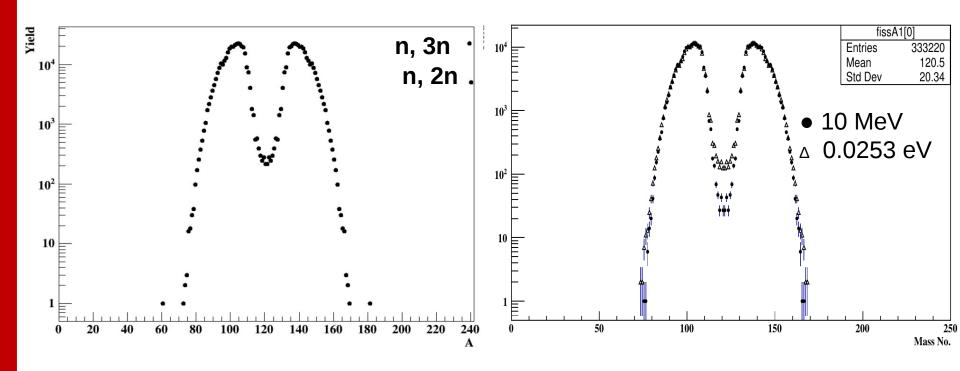
Prompt neutron spectrum is harder than delayed neutron spectrum





Sampling: PU-241 Fission fragments

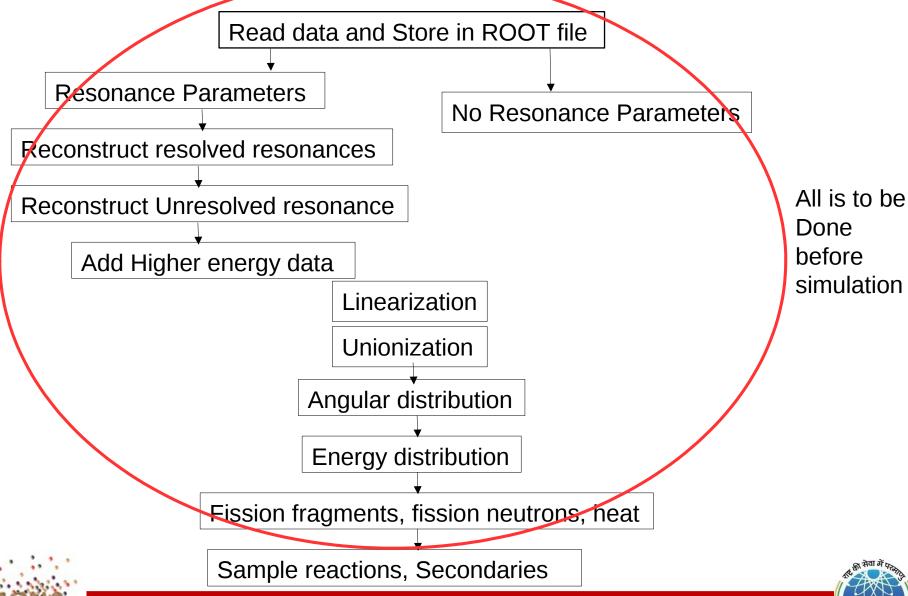
Data are given mostly for 2-3 energies (0.0253 eV, 0.5MeV, 14MeV). Interpolate for intermediate intervals.







Reconstruction and sampling



Summary

- ➤ Neutron cross-sections are reconstructed and agreement with NJOY data is good.
- > Angle and energy distributions are well described.
- Fission fragments and fission neutron multiplicity are validated.
- ➤ Doppler broadening at various temperature is too many sets of data points (every 50Kelvin)
- ➤One should make effort to parameterize the data for temperature dependence and use them for given temperature (300k-3000K).





Future work

- Transport and Validation for the multiplicity, distributions
- ➤ Validation of photon distributions
- Include thermal scattering data
- ► Atomic relaxation data
- ➤ Optimize sampling technique
- Include variance reduction techniques
- ➤ Optimize for vectorized architecture







Thank you for your attention!



